## SPIN STRUCTURES AND SPECTRA OF $\mathbb{Z}_2^k$ -MANIFOLDS.

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ABSTRACT. We give necessary and sufficient conditions for the existence of  $\operatorname{pin}^{\pm}$  and spin structures on Riemannian manifolds with holonomy group  $\mathbb{Z}_2^k$ . For any  $n \geq 4$  (resp.  $n \geq 6$ ) we give examples of pairs of compact manifolds (resp. compact orientable manifolds)  $M_1$ ,  $M_2$ , non homeomorphic to each other, that are Laplace isospectral on functions and on p-forms for any p and such that  $M_1$  admits a  $\operatorname{pin}^{\pm}$  (resp. spin) structure whereas  $M_2$  does not.

### Introduction

Any Riemannian manifold M has naturally associated differential operators of second order, the Laplacian  $\Delta$  acting on smooth functions and more generally, the p-Laplacian  $\Delta_p$  acting on smooth p-forms for  $0 \le p \le n$ . The Dirac operator D is a first order operator that can not always be defined. To make this possible, M needs to have an additional structure: a spin structure, if M is orientable, and a pin<sup>±</sup> structure, in general. In this case one says that M is spin or pin<sup>±</sup>, respectively.

In this paper we consider a question posed by David Webb, namely, can one hear the property of being spin on a compact Riemannian manifold? We shall answer this question in the negative by giving several examples of Laplace isospectral Riemannian manifolds  $M_1, M_2$  such that  $M_1$  is spin (resp.  $\operatorname{pin}^{\pm}$ ) but  $M_2$  has no spin (resp.  $\operatorname{pin}^{\pm}$ ) structure. All our examples will be isospectral on p-forms for  $0 \leq p \leq n$  and will be given by  $\mathbb{Z}_2^k$ -manifolds, that is, compact Riemannian manifolds with holonomy group  $F \simeq \mathbb{Z}_2^k$ . We note that by the Cartan-Ambrose-Singer theorem, such a manifold is necessarily flat, hence of the form  $M_{\Gamma} = \Gamma \backslash \mathbb{R}^n$ ,  $\Gamma$  a Bieberbach group.

In one of the main results, Theorem 2.1, we give a parametrization of the  $\operatorname{pin}^{\pm}$  or spin structures of  $M_{\Gamma}$ , showing that the number is either  $2^r$  for some  $r \geq k$  or zero, and deriving a simple criterion for non existence (see Remark 2.3). In Section 3 we apply Theorem 2.1 and this criterion to construct several isospectral pairs M, M' of  $\mathbb{Z}_2^2$ -manifolds of dimensions  $n \geq 4$  (resp.  $n \geq 6$ ), such that M admits a  $\operatorname{pin}^{\pm}$  (resp. spin) structure while M' does not, thus giving a negative answer to Webb's question. By increasing dimensions, we obtain examples of pairs having these same properties and with the extra condition that both M, M' are Kähler (see Remark 3.1).

In the last section we specialize to the case k=1, i.e. of  $\mathbb{Z}_2$ -manifolds. We show that any such  $M_{\Gamma}$  has  $2^{n-j}$  pin<sup>±</sup> structures for some  $0 \le j \le \lfloor \frac{n-1}{2} \rfloor$ ,

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with j determined by the  $\mathbb{Z}_2$ -action. If furthermore  $M_{\Gamma}$  is of the so called diagonal type and orientable, it turns out that  $M_{\Gamma}$  admits  $2^n$  spin structures, as in the case of the n-torus (see [Fr]).

#### 1. Preliminaries

Bieberbach manifolds. A crystallographic group is a discrete, cocompact subgroup  $\Gamma$  of the isometry group  $I(\mathbb{R}^n)$  of  $\mathbb{R}^n$ . If  $\Gamma$  is torsion-free, then  $\Gamma$  is said to be a Bieberbach group. Such a  $\Gamma$  acts properly discontinuously on  $\mathbb{R}^n$ , thus  $M_{\Gamma} = \Gamma \backslash \mathbb{R}^n$  is a compact flat Riemannian manifold with fundamental group  $\Gamma$  and furthermore, any such manifold arises in this way. Since  $I(\mathbb{R}^n) \simeq \mathrm{O}(n) \ltimes \mathbb{R}^n$ , any element  $\gamma \in I(\mathbb{R}^n)$  decomposes uniquely as  $\gamma = BL_b$ , with  $B \in \mathrm{O}(n)$  and  $b \in \mathbb{R}^n$ . The translations in  $\Gamma$  form a normal maximal abelian subgroup of finite index  $L_{\Lambda}$ ,  $\Lambda$  a lattice in  $\mathbb{R}^n$  which is B-stable for every  $BL_b \in \Gamma$ . The restriction to  $\Gamma$  of the canonical projection  $r: I(\mathbb{R}^n) \to \mathrm{O}(n)$  given by  $BL_b \mapsto B$  is a homomorphism with kernel  $L_{\Lambda}$  and  $r(\Gamma)$  is a finite subgroup of  $\mathrm{O}(n)$  isomorphic to  $F:=L_{\Lambda}\backslash\Gamma$ . It is called the holonomy group of  $\Gamma$  and gives the linear holonomy group of the Riemannian manifold  $M_{\Gamma}$ .

A Bieberbach group  $\Gamma$  is said to be of diagonal type (see [MR2], Definition 1.3) if there exists an orthonormal  $\mathbb{Z}$ -basis  $\{\lambda_1, \ldots, \lambda_n\}$  of the lattice  $\Lambda$  such that for any element  $BL_b \in \Gamma$ ,  $B\lambda_i = \pm \lambda_i$  for  $1 \leq i \leq n$ . These Bieberbach groups have a rather simple holonomy action, among those with holonomy group  $\mathbb{Z}_2^k$ . If  $\Gamma$  is of diagonal type, after conjugation of  $\Gamma$  by an isometry, it may be assumed that  $\Lambda$  is the canonical lattice and that b lies in  $\frac{1}{2}\Lambda$  for any  $\gamma = BL_b \in \Gamma$ . Thus, any  $\gamma \in \Gamma$  can be written uniquely as  $\gamma = BL_{bo}L_{\lambda}$ , where the coordinates of  $b_o$  are 0 or  $\frac{1}{2}$  and  $\lambda \in \Lambda$  (see [MR2], Lemma 1.4).

Pin and spin groups. For a discussion of the material in this subsection we refer to [LM], [Fr2] or [GLP]. Let  $Cl^{\pm}(n)$  denote the Clifford algebras of  $\mathbb{R}^n$  endowed with the definite quadratic forms  $\mp ||\cdot||^2$ . If  $\{e_1,\ldots,e_n\}$  denotes the canonical basis of  $\mathbb{R}^n$ , then a basis for  $Cl^{\pm}(n)$  is given by the set  $\{e_{i_1}\ldots e_{i_k}: 1\leq i_1<\cdots< i_k\leq n\}$ . On  $Cl^{\pm}(n)$  one has the relation  $vw+wv=\pm 2\langle v,w\rangle$  for any  $v,w\in\mathbb{R}^n$ , where  $\langle , \rangle$  denotes the standard inner product. Thus,

(1.1) 
$$e_i e_j = -e_j e_i \qquad \text{for } i \neq j \text{ for both } Cl^{\pm}(n),$$
$$e_i^2 = \pm 1 \qquad \text{for } 1 \leq i \leq n \text{ in } Cl^{\pm}(n).$$

We have compact Lie subgroups,  $\operatorname{Pin}^{\pm}(n)$ , of the group of units of  $Cl^{\pm}(n)$ , with  $\operatorname{Pin}^{\pm}(n) = \{v_1 \dots v_h : v_j \in \mathbb{R}^n, \|v_j\| = 1, \ 1 \leq j \leq h\}$ . The connected component of the identity in both cases is isomorphic to  $\operatorname{Spin}(n) = \{v_1 \dots v_h : v_j \in \mathbb{R}^n, \|v_j\| = 1, \ 1 \leq j \leq h, h \text{ even}\}$ , a compact, simply connected Lie group for  $n \geq 3$ .

Let  $\alpha$  be the canonical involution of  $Cl^{\pm}(n)$  given by  $\alpha(v_1 \dots v_h) = (-1)^h v_1 \dots v_h$ . Then, we have Lie group epimorphisms

$$\mu_{\pm}: \operatorname{Pin}^{\pm}(n) \to \operatorname{O}(n)$$

with kernel  $\{\pm 1\}$ , given by  $\mu_{\pm}(v)(x) = \alpha(v)xv^{-1}$  where  $v \in \operatorname{Pin}^{\pm}(n)$  and  $x \in \mathbb{R}^n$ . If  $v \in \mathbb{R}^n$ , ||v|| = 1, then  $\mu_{\pm}(v)(x) = -vxv^{-1} = \rho_v(x)$  where  $\rho_v$  denotes the orthogonal reflection with respect to the hyperplane orthogonal

to v. When restricted to the connected component of the identity,  $\mu := \mu_{\pm} : \operatorname{Spin}(n) \simeq \operatorname{Pin}^{\pm}(n)_o \to \operatorname{SO}(n)$  give double coverings.

If  $A_j$  is a matrix, for  $1 \leq j \leq m$  we will abuse notation by denoting by  $\operatorname{diag}(A_1, \ldots, A_m)$  the matrix having  $A_j$  in the "diagonal" position j.

Let  $B(t) = \begin{bmatrix} \cos t - \sin t \\ \sin t & \cos t \end{bmatrix}$  with  $t \in \mathbb{R}$  and put

$$\tau(t_1, \dots, t_m) = \begin{cases} \operatorname{diag}(B(t_1), \dots, B(t_m)), & \text{if } n = 2m \\ \operatorname{diag}(B(t_1), \dots, B(t_m), 1), & \text{if } n = 2m + 1. \end{cases}$$

We have that  $T = \{\tau(t_1, \dots, t_m) : t_j \in \mathbb{R}\}$  is a maximal torus of SO(n). A maximal torus of Spin(n) is given by

$$\tilde{T} = \Big\{ \prod_{j=1}^{m} \left( \cos t_j + \sin t_j \ e_{2j-1} e_{2j} \right) : t_j \in \mathbb{R} \Big\}.$$

The restriction  $\mu: \tilde{T} \to T$  is a 2-fold cover and

(1.2) 
$$\mu\left(\prod_{j=1}^{m}(\cos t_j + \sin t_j \ e_{2j-1}e_{2j})\right) = \tau(2t_1, \dots, 2t_m).$$

Spin structures and  $pin^{\pm}$  structures. If (M,g) is a Riemannian manifold of dimension n, let  $B(M) = \bigcup_{x \in M} B_x(M)$  be the bundle of frames on M and  $\pi$ :  $B(M) \to M$  the canonical projection. That is, for  $x \in M$ ,  $B_x(M)$  is the set of ordered orthonormal bases  $(v_1, \ldots, v_n)$  of  $T_x(M)$  and  $\pi((v_1, \ldots, v_n)) = x$ . B(M) is a principal O(n)-bundle over M and, if M is orientable, the bundle of oriented frames  $B^+(M)$  is a principal SO(n)-bundle. A  $pin^{\pm}$  structure on M is a 2-fold cover  $p: \tilde{B}(M) \to B(M)$  that is equivariant and so that  $\tilde{\pi}: \tilde{B}(M) \to M$  is a principal  $Pin^{\pm}(n)$ -bundle with  $\pi \circ p = \tilde{\pi}$ . Similarly, a spin structure on an orientable manifold M is an equivariant 2-fold cover  $p: \tilde{B}^+(M) \to B^+(M)$  where  $\tilde{\pi}: \tilde{B}^+(M) \to M$  is a principal Spin(n)-bundle and  $\pi \circ p = \tilde{\pi}$ .

A manifold in which a spin or a  $pin^{\pm}$  structure has been chosen is called a spin or a  $pin^{\pm}$  manifold, respectively. Note that if M is orientable, any  $pin^{\pm}$  structure on M defines a spin structure and conversely.

We will be interested on spin and pin<sup>±</sup> structures on quotients  $M_{\Gamma} = \Gamma \backslash \mathbb{R}^n$ , where  $\Gamma$  is a Bieberbach group. If  $M = \mathbb{R}^n$ , we have that  $B(\mathbb{R}^n) = \mathbb{R}^n \times O(n)$ , thus clearly  $\mathbb{R}^n \times \operatorname{Pin}^{\pm}(n)$  are principal  $\operatorname{Pin}^{\pm}(n)$ -bundles and the maps  $Id \times \mu_{\pm} : \mathbb{R}^n \times \operatorname{Pin}^{\pm}(n) \to \mathbb{R}^n \times \operatorname{O}(n)$  are equivariant 2-fold covering maps. Similarly, we have that  $\mathbb{R}^n \times \operatorname{Spin}(n)$  is a principal  $\operatorname{Spin}(n)$ -bundle and an equivariant 2-fold cover of  $B^+(\mathbb{R}^n) = \mathbb{R}^n \times \operatorname{SO}(n)$ . Thus we have spin and  $\operatorname{pin}^{\pm}$  structures on  $\mathbb{R}^n$  and since  $\mathbb{R}^n$  is contractible these are the only such structures. Now, if  $\Gamma$  is a Bieberbach group we have a left action of  $\Gamma$  on B(M) given by  $\gamma \cdot (x, (w_1, \dots, w_n)) = (\gamma x, (\gamma_* w_1, \dots, \gamma_* w_n))$ . If  $\gamma = BL_b$  then  $\gamma_* w_j = w_j B$ . Fix  $(v_1, \dots, v_n) \in B(M)$ . Since  $(w_1, \dots, w_n) = (v_1 k, \dots, v_n k)$  for some  $k \in O(n)$ , we see that  $\gamma_* w_j = (v_j k)B = v_j(Bk)$ , thus the action of  $\Gamma$  on B(M) corresponds to the action of  $\Gamma$  on  $\mathbb{R}^n \times O(n)$  given by  $\gamma \cdot (x, k) = (\gamma x, Bk)$ .

Now assume that there is a group homomorphism  $\varepsilon : \Gamma \to \operatorname{Spin}(n)$  (resp.  $\varepsilon_{\pm} : \Gamma \to \operatorname{Pin}^{\pm}(n)$ ) such that  $\mu(\varepsilon(\gamma)) = r(\gamma)$  (resp.  $\mu_{\pm}(\varepsilon_{\pm}(\gamma)) = r(\gamma)$ ). In this case we can lift the left action of  $\Gamma$  on  $B^{+}(\mathbb{R}^{n})$  (resp. on  $B(\mathbb{R}^{n})$ ) to

 $\tilde{\operatorname{B}}^+(\mathbb{R}^n) = \mathbb{R}^n \times \operatorname{Spin}(n)$  (resp. to  $\tilde{\operatorname{B}}(\mathbb{R}^n) = \mathbb{R}^n \times \operatorname{Pin}^{\pm}(n)$ ) via  $\gamma \cdot (x, \tilde{k}) = (\gamma x, \varepsilon(\gamma)\tilde{k})$ . Thus we have the spin structure

$$\Gamma\backslash(\mathbb{R}^n\times\mathrm{Spin}(n))\xrightarrow{\overline{\mathrm{Id}\times\mu}}\Gamma\backslash(\mathbb{R}^n\times\mathrm{SO}(n))$$

$$\Gamma\backslash\mathbb{R}^n$$

for  $M_{\Gamma}$  since  $\Gamma \backslash B(\mathbb{R}^n) = B(\Gamma \backslash \mathbb{R}^n)$  and  $\overline{Id \times \mu}$  is equivariant. Similarly for the pin<sup>±</sup> structures.

In this way, for each homomorphism  $\varepsilon$  or  $\varepsilon_{\pm}$  as above, we obtain a spin or a pin<sup>±</sup> structure on  $M_{\Gamma}$ , respectively. It turns out that all spin and pin<sup>±</sup> structures on  $M_{\Gamma}$  are obtained in this manner (see [**Fr2**], [**LM**]).

The *n*-torus admits  $2^n$  spin structures. Indeed, if  $T_{\Lambda} = \Lambda \backslash \mathbb{R}^n$ , and  $\lambda_1, \ldots, \lambda_n$  is a  $\mathbb{Z}$ -basis of  $\Lambda$ , then a homomorphism  $\varepsilon$  as above is determined by the *n*-tuple  $\varepsilon(L_{\lambda_i}) = \delta_i \in \{\pm 1\}$ , for  $1 \leq i \leq n$  (see [**Fr**]). We shall show in Section 4 that this is still the number of such structures for flat manifolds with holonomy group  $\mathbb{Z}_2$  which are of diagonal type.

# 2. Spin and $\operatorname{pin}^{\pm}$ structures on $\mathbb{Z}_2^k$ -manifolds.

In this section we study the existence of  $pin^{\pm}$  structures on  $\mathbb{Z}_2^k$ -manifolds, showing that the number of such structures is either 0 or  $2^r$  for some  $r \geq k$ . As an application, in the next section we will construct many examples of  $\mathbb{Z}_2^2$ -manifolds for any  $n \geq 4$ , having  $pin^+$  structures but no  $pin^-$  structures (and conversely) or else, having neither of them.

Let  $\Gamma$  be a Bieberbach group with holonomy group  $F \simeq \mathbb{Z}_2^k$ ,  $1 \leq k \leq n-1$ , and translation lattice  $\Lambda$ . Then  $M_{\Gamma} = \Gamma \backslash \mathbb{R}^n$  with  $\Gamma = \langle \gamma_1, \ldots, \gamma_k, \Lambda \rangle$  where  $\gamma_i = B_i L_{b_i}$ ,  $B_i \in O(n)$ ,  $b_i \in \mathbb{R}^n$ ,  $B_i \Lambda = \Lambda$ ,  $B_i^2 = Id$  and  $B_i B_j = B_j B_i$ , for each  $1 \leq i, j \leq k$ .

Assume there is a pin<sup>±</sup> structure on  $M_{\Gamma}$ , that is, a group homomorphism  $\varepsilon_{\pm}: \Gamma \to \operatorname{Pin}^{\pm}(n)$  such that  $\mu_{\pm} \circ \varepsilon_{\pm} = r$ . Then, necessarily  $\varepsilon_{\pm}(L_{\lambda}) \in \{\pm 1\}$ , for  $\lambda \in \Lambda$ . Thus, if  $\lambda_1, \ldots, \lambda_n$  is a  $\mathbb{Z}$ -basis of  $\Lambda$  and we set  $\delta_i := \varepsilon_{\pm}(L_{\lambda_i})$ , for every  $\lambda = \sum_i m_i \lambda_i \in \Lambda$  with  $m_i \in \mathbb{Z}$ , we have  $\varepsilon_{\pm}(L_{\lambda}) = \prod_i \delta_i^{m_i} = \prod_{m_i \text{ odd}} \delta_i$ . If  $\gamma = RL_i \in \Gamma$  we will fix a distinguished (though arbitrary) element in

If  $\gamma = BL_b \in \Gamma$  we will fix a distinguished (though arbitrary) element in  $\mu_{\pm}^{-1}(B)$ , denoted by  $u_{\pm}(B)$ . If  $M_{\Gamma}$  is orientable, we write  $u(B) := u_{\pm}(B)$ . Thus, if  $\gamma = BL_b \in \Gamma$ , then

(2.1) 
$$\varepsilon_{\pm}(\gamma) = \sigma \, u_{\pm}(B),$$

where  $\sigma \in \{\pm 1\}$  depends on  $\gamma$  and on the choice of  $u_{\pm}(B)$ .

Let  $\Gamma = \langle \gamma_1, \dots, \gamma_k, \Lambda \rangle$ . The morphism  $\varepsilon_{\pm}$  is determined by its action on the generators of  $\Gamma$ , that is, by the (n+k)-tuple

(2.2) 
$$(\delta_1, \dots, \delta_n, \sigma_1 u_{\pm}(B_1), \dots, \sigma_k u_{\pm}(B_k))$$
 or 
$$(\delta_1, \dots, \delta_n, \sigma_1, \dots, \sigma_k) \in \{\pm 1\}^{n+k}$$

where  $\delta_i = \varepsilon_{\pm}(L_{\lambda_i})$  and  $\sigma_i$  is defined by the equation  $\varepsilon_{\pm}(\gamma_i) = \sigma_i u_{\pm}(B_i)$ , for  $1 \leq i \leq k$ .

Now, since  $\varepsilon_{\pm}$  is a homomorphism, for any  $\gamma = BL_b \in \Gamma, \lambda \in \Lambda$  we have

$$\varepsilon_{\pm}(L_{B\lambda}) = \varepsilon_{\pm}(\gamma L_{\lambda} \gamma^{-1}) = \varepsilon_{\pm}(\gamma) \varepsilon_{\pm}(L_{\lambda}) \varepsilon_{\pm}(\gamma^{-1}) = \varepsilon_{\pm}(L_{\lambda}).$$

Therefore we see that if  $\varepsilon_{\pm}$  is a pin<sup> $\pm$ </sup> structure on  $M_{\Gamma}$ , since  $\gamma^2 \in L_{\Lambda}$ , then the character  $\varepsilon_{\pm|\Lambda}$  must satisfy the following conditions for any  $\gamma = BL_b \in \Gamma$ :

(2.3) 
$$(\varepsilon_1) \qquad \qquad \varepsilon_{\pm}(\gamma^2) = \varepsilon_{\pm}(\gamma)^2 = u_{\pm}^2(B)$$

$$(\varepsilon_2) \qquad \qquad \varepsilon_{\pm}(L_{(B-Id)\lambda}) = 1, \quad \lambda \in \Lambda.$$

We thus set

(2.4) 
$$\hat{\Lambda}(\Gamma) := \{ \chi \in \text{Hom}(\Lambda, \{\pm 1\}) : \chi \text{ satisfies } (\varepsilon_1) \text{ and } (\varepsilon_2) \}.$$

The next result gives a parametrization of the pin<sup>±</sup> structures  $\varepsilon_{\pm}$  for  $M_{\Gamma}$ .

**Theorem 2.1.** If  $\Gamma = \langle \gamma_1, \dots, \gamma_k, \Lambda \rangle$  is a Bieberbach group with holonomy group  $\mathbb{Z}_2^k$  and  $\sigma_1, \dots, \sigma_k$  are as in (2.2), then the map  $\varepsilon_{\pm} \mapsto (\varepsilon_{\pm|\Lambda}, \sigma_1, \dots, \sigma_k)$  defines a bijective correspondence between the pin<sup>±</sup> structures on  $M_{\Gamma}$  and the set  $\hat{\Lambda}(\Gamma) \times \{\pm 1\}^k$ . The number of pin<sup>±</sup> structures on  $M_{\Gamma}$  is either 0 or  $2^r$  for some  $r \geq k$ .

*Proof.* We shall write  $\varepsilon, \mu, u(B)$  in place of  $\varepsilon_{\pm}, \mu_{\pm}, u_{\pm}(B)$ , for simplicity.

Any element  $\gamma \in \Gamma$  can be written as a product of generators  $\gamma_i = B_i L_{b_i}$  and  $\lambda \in \Lambda$ . After reordering, by normality of  $\Lambda$  in  $\Gamma$  and since  $B_i^2 = Id$ , we see that  $\gamma$  can be written uniquely as

(2.5) 
$$\gamma = \gamma_{i_1} \dots \gamma_{i_r} L_{\lambda}, \quad \text{with } 1 \le i_1 < \dots < i_r \le k, \ \lambda \in \Lambda.$$

Given  $\varepsilon \in \hat{\Lambda}(\Gamma)$  and for any choices of  $\varepsilon(\gamma_i) \in \mu^{-1}(B_i)$ ,  $1 \le i \le k$ , we define (in the notation of (2.5)) for  $\gamma \in \Gamma$ :

(2.6) 
$$\varepsilon(\gamma) = \varepsilon(\gamma_{i_1}) \dots \varepsilon(\gamma_{i_r}) \varepsilon(L_{\lambda}).$$

Thus, we get a well defined map  $\varepsilon : \Gamma \to \operatorname{Pin}^{\pm}(n)$  such that  $\mu \circ \varepsilon = r$  and we claim it is a homomorphism. For this purpose we need to show that

(2.7) 
$$\varepsilon(\gamma_{i_1} \dots \gamma_{i_r} L_{\lambda} \gamma_{j_1} \dots \gamma_{j_t} L_{\lambda'}) = \varepsilon(\gamma_{i_1} \dots \gamma_{i_r} L_{\lambda}) \varepsilon(\gamma_{j_1} \dots \gamma_{j_t} L_{\lambda'}),$$

for any  $i_1 < \cdots < i_r$ ,  $j_1 < \cdots < j_t$  and  $\lambda, \lambda' \in \Lambda$ .

We first note that we may leave out  $\lambda, \lambda'$  in (2.7). Indeed, assume that for  $\gamma, \gamma' \in \Gamma$  one has  $\varepsilon(\gamma\gamma') = \varepsilon(\gamma)\varepsilon(\gamma')$ . Then, by  $(\varepsilon_2)$ 

$$\varepsilon(\gamma L_{\lambda} \gamma' L_{\lambda'}) = \varepsilon(\gamma \gamma' L_{B\lambda + \lambda'}) = \varepsilon(\gamma \gamma') \varepsilon(L_{B\lambda + \lambda'}) 
= \varepsilon(\gamma) \varepsilon(\gamma') \varepsilon(L_{\lambda}) \varepsilon(L_{\lambda'}) = \varepsilon(\gamma L_{\lambda}) \varepsilon(\gamma' L_{\lambda'}).$$

As a step in the proof of (2.7) (with  $\lambda = \lambda' = 0$ ) we will first show that

(2.8) 
$$\varepsilon(\gamma_i \gamma_j) = \varepsilon(\gamma_i) \varepsilon(\gamma_j), \text{ for any } i, j.$$

This follows from the definition of  $\varepsilon$ , if i < j, and from condition  $(\varepsilon_1)$ , if i = j. We thus assume that j < i. Then we may write  $\gamma_i \gamma_j = \gamma_j \gamma_i [\gamma_i^{-1}, \gamma_j^{-1}]$ . Since  $[\gamma_i^{-1}, \gamma_j^{-1}] \in \Lambda$ , by the definition of  $\varepsilon$ 

(2.9) 
$$\varepsilon(\gamma_i \gamma_j) = \varepsilon(\gamma_j) \varepsilon(\gamma_i) \varepsilon([\gamma_i^{-1}, \gamma_j^{-1}]).$$

Note that (2.9) will equal  $\varepsilon(\gamma_i)\varepsilon(\gamma_j)$  if and only if it holds the relation

(2.10) 
$$\varepsilon([\gamma_i^{-1}, \gamma_j^{-1}]) = [\varepsilon(\gamma_i^{-1}), \varepsilon(\gamma_j^{-1})].$$

To show (2.10), we have by condition  $(\varepsilon_1)$  that

(2.11) 
$$\varepsilon((\gamma_j \gamma_i)^2) = \varepsilon(\gamma_j \gamma_i)^2 = \varepsilon(\gamma_j)\varepsilon(\gamma_i)\varepsilon(\gamma_j)\varepsilon(\gamma_i).$$

On the other hand

(2.12) 
$$\varepsilon((\gamma_{j}\gamma_{i})^{2}) = \varepsilon(\gamma_{j}\gamma_{j}\gamma_{i}[\gamma_{i}^{-1}, \gamma_{j}^{-1}]\gamma_{i}) \\
= \varepsilon(\gamma_{j}^{2}\gamma_{i}^{2}(\gamma_{i}^{-1}[\gamma_{i}^{-1}, \gamma_{j}^{-1}]\gamma_{i})) \\
= \varepsilon(\gamma_{j}^{2})\varepsilon(\gamma_{i}^{2})\varepsilon(\gamma_{i}^{-1}[\gamma_{i}^{-1}, \gamma_{j}^{-1}]\gamma_{i}) \\
= \varepsilon(\gamma_{j})^{2}\varepsilon(\gamma_{i})^{2}\varepsilon([\gamma_{i}^{-1}, \gamma_{j}^{-1}]).$$

In the last equality we have used condition  $(\varepsilon_2)$  and the fact that commutators lie in  $\Lambda$ .

Now, by combining (2.11) and (2.12) we obtain (2.10), hence (2.8) follows.

In the general case, (2.7) can be proved by an inductive argument.

Let first t=1, r arbitrary. The case r=1 is (2.8), so assume r>1. If  $j_1>i_r$ , then the assertion is clear by the definition of  $\varepsilon$ , while if  $j_1=i_r$ , we may use  $(\varepsilon_1)$  and induction. We thus assume that there is  $\alpha$  such that  $i_{\alpha-1} \leq j_1 < i_{\alpha}$ . Actually, we shall take  $i_{\alpha-1} < j_1$ . The proof when  $i_{\alpha-1} = j_1$  is similar, but simpler.

If we set  $u = \gamma_{i_{\alpha}} \cdots \gamma_{i_r}$  then

$$\varepsilon(\gamma_{i_1}\cdots\gamma_{i_r}\gamma_{j_1}) = \varepsilon(\gamma_{i_1}\cdots\gamma_{i_{\alpha-1}}\gamma_{j_1}u[u^{-1},\gamma_{j_1}^{-1}]) 
= \varepsilon(\gamma_{i_1})\cdots\varepsilon(\gamma_{i_{\alpha-1}})\varepsilon(\gamma_{j_1})\varepsilon(u)\varepsilon([u^{-1},\gamma_{j_1}^{-1}]) 
\text{(by (2.10))} = \varepsilon(\gamma_{i_1})\cdots\varepsilon(\gamma_{i_{\alpha-1}})\varepsilon(u)\varepsilon(\gamma_{j_1}) 
= \varepsilon(\gamma_{i_1}\cdots\gamma_{i_r})\varepsilon(\gamma_{j_1}).$$

The argument for arbitrary t is quite similar and will be omitted.  $\Box$ 

Remark 2.2. For manifolds of diagonal type, condition  $(\varepsilon_2)$  always holds, since  $(B-Id)\Lambda \subset 2\Lambda$  for any  $BL_b \in \Gamma$ . More generally, for manifolds whose holonomy representation decomposes as a sum of integral representations of rank  $\leq 2$ , condition  $(\varepsilon_2)$  can be expressed in simple terms.

In Section 4 we will study in more detail the case of  $\mathbb{Z}_2$ -manifolds, showing in particular that  $pin^{\pm}$  structures can always be defined in this case.

**Remark 2.3.** The previous theorem shows that there are restrictions for a  $\mathbb{Z}_2^k$ -manifold  $M_{\Gamma}$  to carry a pin<sup>±</sup> structure. As a consequence, one has the following simple criterion:

Suppose there exist  $\gamma = BL_b, \gamma' = B'L_{b'} \in \Gamma$  with  $\gamma^2 = {\gamma'}^2$  and such that for  $u_+(B) \in \mu_+^{-1}(B)$  and  $u_+(B') \in \mu_+^{-1}(B')$  one has  $u_+(B)^2 = -u_+(B')^2$ . Then  $M_{\Gamma}$  can not admit a pin<sup>+</sup> structure.

Indeed, such a structure  $\varepsilon_+$  would have to satisfy  $\varepsilon_+(\gamma) = \pm u_+(B)$ ,  $\varepsilon_+(\gamma') = \pm u_+(B')$  and  $\varepsilon_+(\gamma^2) = \varepsilon_+({\gamma'}^2)$ , that is,  $u_+(B)^2 = u_+(B')^2$  against our assumption.

The same criterion, with the obvious changes, is valid for non existence of pin<sup>-</sup> structures, or spin structures in the orientable case.

**Remark 2.4.** In contrast with Remark 2.3, by applying the doubling procedure in  $[\mathbf{DM2}]$ , we may obtain spin Bieberbach manifolds of diagonal type with holonomy group  $\mathbb{Z}_2^k$ , for any  $k \geq 1$ . Indeed, let  $\Gamma = \langle \gamma_1, \ldots, \gamma_k, L_\Lambda \rangle$  be an n-dimensional Bieberbach group of diagonal type with holonomy group  $\mathbb{Z}_2^k$ . Define  $d\Gamma := \langle d\gamma_1, \ldots, d\gamma_k, L_{\Lambda \oplus \Lambda} \rangle$  where  $d\gamma := \begin{bmatrix} B & 0 \\ 0 & B \end{bmatrix} L_{(b,b)}$  if  $\gamma = BL_b \in \Gamma$  (see Definition 3.1 in  $[\mathbf{DM2}]$ ). Thus,  $d\Gamma$  is a Bieberbach group

of dimension 2n with holonomy group  $\mathbb{Z}_2^k$ . The manifold  $M_{\mathrm{d}\Gamma} = \mathrm{d}\Gamma \backslash \mathbb{R}^{2n}$  is an orientable Kähler flat manifold of diagonal type. If we apply this procedure twice, then the manifold  $M_{\mathrm{d}^2\Gamma}$  is hyperkähler (see Proposition 3.2 in  $[\mathbf{DM2}]$ ). It turns out that this 4n-dimensional manifold is always spin. Indeed, in the notation of Lemma 3.1 in the next section, since  $h \in 4\mathbb{Z}$  for  $\mathrm{d}^2\Gamma$ , we have that  $u^2(B) = u_{0,h}^2 = 1$  by (3.3). Hence, condition  $(\varepsilon_1)$  takes the form  $\varepsilon(\gamma^2) = 1$  for any  $\gamma \in \Gamma$ . Therefore, spin structures can always be defined for  $M_{\mathrm{d}^2\Gamma}$ , for example we may take any of the  $2^k$  homomorphisms  $\varepsilon : \Gamma \to \mathrm{Spin}(n)$  such that  $\varepsilon_{|\Lambda} \equiv 1$ .

### 3. Spin structures on some isospectral pairs.

In this section we will construct several isospectral pairs  $\{M, M'\}$  of  $\mathbb{Z}_2^2$ -manifolds of dimension 4 by using the results in [MR2], and we will determine the pin<sup>±</sup> or spin structures, showing that, for some of them, M has a pin<sup>±</sup> or a spin structure, while M' does not. The main result is given in Theorem 3.2. In the proof, we will need to know some preimages in  $\operatorname{Pin}^{\pm}(n)$  by  $\mu_+$ , as well as their squares.

Set  $J := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ . For each  $0 \le j, h < n$ , we set

(3.1) 
$$B_{j,h} = \operatorname{diag}(\underbrace{J, \dots, J}_{j}, \underbrace{-1, \dots, -1}_{h}, \underbrace{1, \dots, 1}_{l}),$$

where n = 2j + h + l,  $j + h \neq 0$  and  $l \geq 1$ .

**Lemma 3.1.** Let  $B_{j,h}$  be as in (3.1) and let  $\mu_{\pm} : Pin^{\pm}(n) \to O(n)$  be the canonical covering maps. If we set

$$(3.2) u_{j,h}^{\pm} := (\frac{\sqrt{2}}{2})^j (e_1 - e_2) \dots (e_{2j-1} - e_{2j}) e_{2j+1} \dots e_{2j+h},$$

then  $\mu_{+}^{-1}(B_{j,h}) = \{\pm u_{j,h}^{+}\}, \ \mu_{-}^{-1}(B_{j,h}) = \{\pm u_{j,h}^{-}\}$  and furthermore

(3.3) 
$$(u_{j,h}^{+})^{2} = (-1)^{jh} (-1)^{\left[\frac{j}{2}\right]} (-1)^{\left[\frac{h}{2}\right]}$$
$$(u_{j,h}^{-})^{2} = (-1)^{jh} (-1)^{\left[\frac{j+1}{2}\right]} (-1)^{\left[\frac{h+1}{2}\right]}.$$

In particular,  $(u_{0,h}^+)^2 = (-1)^{\left[\frac{h}{2}\right]}$  and  $(u_{0,h}^-)^2 = (-1)^{\left[\frac{h+1}{2}\right]}$ . If  $B_{j,h} \in SO(n)$ , i.e. if j+h is even, then  $u_{j,h}^2 = (-1)^{\frac{j+h}{2}}$ .

If  $B \in O(n)$  is conjugate to  $B_{j,h}$ , and  $u_{\pm}(B) \in \mu_{\pm}^{-1}(B)$ , then  $u_{\pm}^{2}(B) = (u_{j,h}^{\pm})^{2}$ .

*Proof.* Since  $\mu_{\pm}(e_i) = \rho_{e_i} = \operatorname{diag}(1, \dots, 1, -1, 1, \dots, 1)$  with -1 in the i-th position, it is clear that  $\mu_+^{-1}(B_{0,h}) = \mu_-^{-1}(B_{0,h}) = \left\{ \pm e_1 \dots e_h \right\}$ . If n = 2, we may write J as a product  $J = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$ . Hence, using (1.2), and  $\mu_{\pm}(e_i) = \rho_{e_i}$ , we get that  $\mu_+^{-1}(J) = \mu_-^{-1}(J) = \left\{ \pm e_1(\cos(\frac{\pi}{4}) + \sin(\frac{\pi}{4})e_1e_2) \right\} = \left\{ \pm \frac{\sqrt{2}}{2}(e_1 - e_2) \right\}$ . Arguing similarly for arbitrary n, the first assertion in the lemma follows.

On the other hand one computes, using (1.1), that both  $(e_1 ldots e_h)^2$  and  $2^{-h}((e_1 - e_2) ldots (e_{2h-1} - e_{2h}))^2$  equal  $(-1)^{\left[\frac{h}{2}\right]}$  in  $Cl^+(n)$  and  $(-1)^{\left[\frac{h+1}{2}\right]}$  in  $Cl^-(n)$ , respectively. This implies equations (3.3).

Now, suppose  $B = CB_{j,h}C^{-1}$  with  $C \in O(n)$ . If  $u_+(C) \in \mu_+(C)^{-1}$ , then  $u_+(B) = \pm u_+(C)u_{j,h}^+u_+(C)^{-1}$  and hence  $u_+^2(B) = u_+(C)(u_{j,h}^+)^2u_+(C)^{-1} = (u_{j,h}^+)^2$ . The verification for  $u_-^2(B)$  is identical.

We now consider some pairs of 4-dimensional  $\mathbb{Z}_2^2$ -manifolds  $\{M_i, M_i'\}$ ,  $1 \leq i \leq 5$ , where  $M_i = \Gamma_i \backslash \mathbb{R}^4$ ,  $M_i' = \Gamma_i' \backslash \mathbb{R}^4$  and the groups  $\Gamma_i = \langle \gamma_1, \gamma_2, \Lambda \rangle$ ,  $\Gamma_i' = \langle \gamma_1', \gamma_2', \Lambda \rangle$  are given in Table 1, where  $\gamma_i = B_i L_{b_i}$ ,  $\gamma_i' = B_i L_{b_i'}$ , i = 1, 2,  $B_3 = B_1 B_2$ ,  $b_3 = B_2 b_1 + b_2$ ,  $b_3' = B_2' b_1' + b_2'$  and  $\Lambda = \mathbb{Z} e_1 \oplus \ldots \oplus \mathbb{Z} e_n$  is the canonical lattice. Furthermore, we take  $B_i = B_i'$ . In all cases the matrices  $B_i$  are diagonal and are written as column vectors. We indicate the translation vectors  $b_i, b_i'$  also as column vectors, leaving out the coordinates that are equal to zero. We will also use the pair  $\{\tilde{M}_1, \tilde{M}_1'\}$  of  $\mathbb{Z}_2^2$ -manifolds of dimension 6 obtained from the pair  $\{M_1, M_1'\}$  by adjoining the characters (-1, 1, -1) and (1, -1, -1) to  $B_i$ ,  $1 \leq i \leq 3$ , and keeping  $b_i, b_i'$  unchanged.

Table 1

(M M/)	$B_1$	$L_{b_1}$	$L_{b'_1}$	$B_2$	$L_{b_2}$	$L_{b_2'}$	$B_3$	$L_{b_3}$	$L_{b_3'}$
	1			1	1/2	1/2	1	1/2	1/2
	1		1/2	1	1/2		1	1/2	1/2
$\{M_1, M_1'\} $ $\{\tilde{M}_1, \tilde{M}_1'\}$	1	1/2		-1			-1	1/2	
$\{M_1,M_1'\}$	-1			1		1/2	-1		1/2
	-1			1			-1		
	1			-1			-1		

	$B_1$	$L_{b_1}$	$L_{b'_1}$	$B_2$	$L_{b_2}$	$L_{b_2'}$	$B_3$	$L_{b_3}$	$L_{b_3'}$
	1			1		1/2	1		1/2
$\{M_2, M_2'\}$	1		1/2	1	1/2	1/2	1	1/2	
	1	1/2		-1			-1	1/2	
	-1			1	1/2		-1	1/2	

	$B_1$	$L_{b_1}$	$L_{b_1'}$	$B_2$	$L_{b_2}$	$L_{b_2'}$	$B_3$	$L_{b_3}$	$L_{b_3'}$
	1			-1			-1		
$\{M_3,M_3'\}$	1		1/2	-1			-1		1/2
	-1			-1	1/2		1	1/2	
	1	1/2		1	1/2	1/2	1		1/2

	$B_1$	$L_{b_1}$	$L_{b_1'}$	$B_2$	$L_{b_2}$	$L_{b_2'}$	$B_3$	$L_{b_3}$	$L_{b_3'}$
	1	1/2		-1			-1	1/2	
$\{M_4, M_4'\}$	1	1/2	1/2	-1			-1	1/2	1/2
	-1			-1		1/2	1		1/2
	1		1/2	1	1/2	1/2	1	1/2	

$$\{M_5, M_5'\}$$

$$\begin{bmatrix} B_1 & L_{b_1} & L_{b_1'} & B_2 & L_{b_2} & L_{b_2'} & B_3 & L_{b_3} & L_{b_3'} \\ -1 & & & 1 & 1/2 & -1 & 1/2 \\ -1 & & & -1 & 1/2 & 1/2 & 1 & 1/2 & 1/2 \\ 1 & & 1/2 & -1 & & -1 & 1/2 \\ 1 & 1/2 & & 1 & 1/2 & & 1 \end{bmatrix}$$

We observe that only  $M_5, M'_5, \tilde{M}_1, \tilde{M}'_1$  are orientable.

In order to show the isospectrality of these pairs we will need to recall some known results.

For  $BL_b \in \Gamma$  set  $n_B := \dim(\mathbb{R}^n)^B = |\{1 \le i \le n : Be_i = e_i\}|$  and

(3.4) 
$$n_B(\frac{1}{2}) := |\{1 \le i \le n : Be_i = e_i \text{ and } b \cdot e_i = \frac{1}{2}\}|.$$

If  $0 \le t \le d \le n$ , the Sunada numbers for  $\Gamma$  are defined by

(3.5) 
$$c_{d,t}(\Gamma) := \left| \left\{ BL_b \in \Gamma : n_B = d \text{ and } n_B(\frac{1}{2}) = t \right\} \right|.$$

In [MR2], Theorem 3.3, it is shown that the equality of the Sunada numbers  $c_{d,t}(\Gamma) = c_{d,t}(\Gamma')$  for every d,t, is equivalent to the validity of the conditions in Sunada's theorem (see [Su]) for  $M_{\Gamma}$  and  $M_{\Gamma'}$ . In particular this implies that  $M_{\Gamma}$  and  $M_{\Gamma'}$  are isospectral on p-forms for  $0 \le p \le n$ . This method was used in [MR] and [MR3] to prove the isospectrality of the pairs  $M_5, M_5'$  and  $M_2, M_2'$  respectively. Also, the method of adding characters and keeping isospectrality was also used in [MR].

We are now in a position to state the main result in this paper.

**Theorem 3.2.** The pairs  $M_i, M'_i, 1 \le i \le 5$ , and  $\tilde{M}_1, \tilde{M}'_1$  are pairwise isospectral.

The number of  $pin^{\pm}$  and spin structures on  $M_i, M'_i, 1 \leq i \leq 5$ , and  $\tilde{M}_1, \tilde{M}'_1$  are given in the following table.

Pairs	$M_1$	$M'_1$	$\tilde{M}_1$	$\tilde{M}_1'$	$M_2$	$M_2'$	$M_3$	$M_3'$	$M_4$	$M'_4$	$M_5$	$M_5'$
$pin^+$		$2^3$		$2^5$	$2^3$	_	_	$2^{4}$	$2^4$	$2^3$	$2^{4}$	$2^3$
pin-	$2^4$	$2^3$	_	$2^5$	$2^3$	-	_	-	_	$2^3$	$2^4$	$2^3$
spin	_	_	_	$2^5$	_	_	_	_	_	_	$2^{4}$	$2^3$

The various isospectral pairs in the table show that one can not hear the existence of  $pin^{\pm}$  or spin structures on a compact Riemannian manifold.

Proof. Since all manifolds are of diagonal type, to show that these pairs are isospectral it suffices to check the equality of the Sunada numbers (see (3.5)). It is easy to see from Table 1 that the non trivial Sunada numbers, besides  $c_{4,0}=1$  corresponding to the identity, are:  $c_{2,2}=c_{3,1}=c_{3,2}=1$  for  $M_1$  and  $M'_1$ ;  $c_{2,2}=c_{4,1}=c_{4,2}=1$  for  $\tilde{M}_1$  and  $\tilde{M}'_1$ ;  $c_{2,1}=c_{3,1}=c_{3,2}=1$  for  $M_2$  and  $M'_2$ ;  $c_{1,1}=c_{2,1}=c_{3,1}=1$  for  $M_3$  and  $M'_3$ ;  $c_{1,1}=c_{2,1}=c_{3,2}=1$  for  $M_4$  and  $M'_4$ ; and  $c_{2,1}=3$  for  $M_5$  and  $M'_5$ . Thus, it follows that all pairs  $M_i, M'_i, 1 \leq i \leq 5$ , and  $\tilde{M}_1, \tilde{M}'_1$  are isospectral on functions.

We shall now use Theorem 2.1 to determine the spin and pin<sup>±</sup> structures on  $M_1, M'_1, \ldots, M_5, M'_5, \tilde{M}_1$  and  $\tilde{M}'_1$ . By Remark 2.2 we need only look at condition  $(\varepsilon_1)$ .

We first look at the pair  $M_1, M'_1$ . We have that

$$\gamma_1^2 = L_{e_3}, \ \gamma_2^2 = L_{e_1+e_2} = \gamma_3^2; \qquad {\gamma_1'}^2 = L_{e_2}, \ {\gamma_2'}^2 = L_{e_1+e_4}, \ {\gamma_3'}^2 = L_{e_1+e_2}.$$
 By (1.2) and Lemma 3.1:

$$u_{\pm}^{2}(B_{1}) = u_{\pm}^{2}(B'_{1}) = (\sigma_{1}e_{4})^{2} = \pm 1, \qquad u_{\pm}^{2}(B_{2}) = u_{\pm}^{2}(B'_{2}) = (\sigma_{2}e_{3})^{2} = \pm 1,$$
  
$$u_{\pm}^{2}(B_{3}) = u_{\pm}^{2}(B'_{3}) = (\sigma_{3}e_{3}e_{4})^{2} = -1$$

with  $\sigma_i \in \{\pm 1\}$ . By the criterion in Remark 2.3, it follows that  $M_1$  has no pin<sup>+</sup> structures, since  $\gamma_2^2 = \gamma_3^2$  and  $u_+^2(B_2) = 1$  while  $u_+^2(B_3) = -1$ .

Furthermore, by the previous equations, if  $\delta_i = \varepsilon_{\pm}(L_{e_i})$ , condition  $(\varepsilon_1)$  gives  $\delta_3 = \pm 1$ ,  $\delta_1 \delta_2 = \pm 1$  and  $\delta_1 \delta_2 = -1$ . The last two equations are not compatible for  $Cl^+(n)$ , hence we see again that  $M_1$  does not admit pin<sup>+</sup> structures. However, it has  $2^4$  pin<sup>-</sup> structures given by

$$\varepsilon_{-}(M_1) = (\delta_1, -\delta_1, -1, \delta_4; \sigma_1 e_4, \sigma_2 e_3)$$

where  $\delta_i, \sigma_j \in \{\pm 1\}$  are arbitrary for i = 1, 4, j = 1, 2. Similarly, condition  $(\varepsilon_1)$  for  $M'_1$  gives  $\delta_2 = \pm 1, \ \delta_1 \delta_4 = \pm \ \text{and} \ \delta_1 \delta_2 = -1$ . Thus,  $M'_1$  has  $2^3 \ \text{pin}^{\pm}$  structures given by

$$\varepsilon_{\pm}(M_1') = (\mp 1, \pm 1, \delta_3, -1; \sigma_1 e_4, \sigma_2 e_3)$$

with  $\delta_3, \sigma_1, \sigma_2 \in \{\pm 1\}$ . In this way, we have shown that  $M_1, M_1'$  is an isospectral pair such that  $M_1$  carries no pin<sup>+</sup> structure while  $M_1'$  admits  $2^3$  of them.

We note that the orientable manifolds  $\tilde{M}_1$ ,  $\tilde{M}'_1$  do have the same properties. These manifolds are still isospectral (again we have equality of Sunada numbers) and  $\gamma_i^2$  and  ${\gamma'_i}^2$  are the same as before, for  $1 \le i \le 3$ .

Now, if we look for spin structures  $\varepsilon$  on  $\tilde{M}_1, \tilde{M}'_1$ , we get

$$u^{2}(B_{1}) = u^{2}(B'_{1}) = (\sigma_{1}e_{4}e_{5})^{2} = -1, \quad u^{2}(B_{2}) = u^{2}(B'_{2}) = (\sigma_{2}e_{3}e_{6})^{2} = -1,$$
  
 $u^{2}(B_{3}) = u^{2}(B'_{3}) = (\sigma_{3}e_{3}e_{4}e_{5}e_{6})^{2} = 1.$ 

For  $\tilde{M}_1$  we have  $\gamma_2^2 = \gamma_3^2 = L_{e_1+e_2}$ , hence  $\varepsilon(\gamma_2^2) = \varepsilon(\gamma_3^2)$ , a contradiction, given that  $u^2(B_2) = -1$  and  $u^2(B_3) = 1$ . Thus, there are no spin structures on  $\tilde{M}_1$ . On the other hand, for  $\tilde{M}_1'$ , we have  ${\gamma_1'}^2 = L_{e_2}$ ,  ${\gamma_2'}^2 = L_{e_1+e_4}$ ,  ${\gamma_3'}^2 = L_{e_1+e_2}$ . Thus,  $\varepsilon(L_{e_2}) = -1$ ,  $\varepsilon(L_{e_1+e_4}) = -1$ ,  $\varepsilon(L_{e_1+e_2}) = 1$ , hence there are  $2^5$  spin structures given by

$$\varepsilon = (-1, -1, \delta_3, 1, \delta_5, \delta_6; \sigma_1 e_4 e_5, \sigma_2 e_3 e_6)$$

with  $\delta_3, \delta_5, \delta_6, \sigma_1, \sigma_2 \in \{\pm 1\}$ .

This proves the claim and shows that one can not hear the existence of spin structures on a compact Riemannian manifold.

We consider next the remaining pairs  $M_i, M'_i, 2 \le i \le 5$ . The calculations are entirely similar to those in the cases discussed above, so we will omit the details, giving the necessary information in several tables. For convenience, we will also include the pair  $M_1, M'_1$ .

Note that the manifolds  $M_1, M_1', M_2, M_2'$ , as well as  $M_3, M_3', M_4, M_4'$ , have the same holonomy representation. Furthermore, all matrices appearing in Table 1 are conjugate to  $B_{0,1}, B_{0,2}$  or  $B_{0,3}$ . By Lemma 3.1 we know that  $u_{0,1}^{\pm 2} = \pm 1, u_{0,2}^{\pm 2} = -1$  and  $u_{0,3}^{\pm 2} = \mp 1$  for  $\operatorname{Pin}^{\pm}(n)$ . Thus we have:

Table 2.

manifolds	$u_{\pm}^{2}(B_{1})$	$u_{\pm}^2(B_2)$	$u_{\pm}^{2}(B_{3})$
$M_1, M_1', M_2, M_2'$	±1	±1	-1
$M_3, M_3', M_4, M_4'$	±1	∓1	-1
$M_5, M'_5$	-1	-1	-1

One has that  $\gamma_i^2 = L_{\lambda_i} \in \Lambda$ . In Table 3 we give the vectors  $\lambda_i$  for  $1 \le i \le 3$  and for every  $M_j, M'_j, 1 \le j \le 5$ .

TABLE	3.

	$M_1$	$M_1'$	$M_2$	$M_2'$	$M_3$	$M_3'$	$M_4$	$M_4'$	$M_5$	$M_5'$
$\gamma_1^2$	$e_3$	$e_2$	$e_3$	$e_2$	$e_4$	$e_2$	$e_1 + e_2$	$e_2 + e_4$	$e_4$	$e_3$
$\gamma_2^2$	$e_1 + e_2$	$e_1 + e_4$	$e_2 + e_4$	$e_1 + e_2$	$e_4$	$e_4$	$e_4$	$e_4$	$e_4$	$e_1$
$\gamma_3^2$	$e_1 + e_2$	$e_1 + e_2$	$e_2$	$e_1$	$e_3$	$e_4$	$e_4$	$e_3$	$e_2$	$e_2$

Using the information obtained in Tables 2 and 3 we get the equations to be satisfied by the  $\delta_i$ 's, resulting from condition ( $\varepsilon_1$ ).

Table 4. Equations for  $\delta_i$ ,  $1 \le i \le 4$ .

	$\gamma_1$	$\gamma_2$	$\gamma_3$
$M_1$	$\delta_3 = \pm 1$	$\delta_1 \delta_2 = \pm 1$	$\delta_1 \delta_2 = -1$
$M'_1$	$\delta_2 = \pm 1$	$\delta_1 \delta_4 = \pm 1$	$\delta_1 \delta_2 = -1$
$M_2$	$\delta_3 = \pm 1$	$\delta_2 \delta_4 = \pm 1$	$\delta_2 = -1$
$M_2'$	$\delta_2 = \pm 1$	$\delta_1 \delta_2 = \pm 1$	$\delta_1 = -1$
$M_3$	$\delta_4 = \pm 1$	$\delta_4 = \mp 1$	$\delta_3 = -1$
$M_3'$	$\delta_2 = \pm 1$	$\delta_4 = \mp 1$	$\delta_4 = -1$
$M_4$	$\delta_1 \delta_2 = \pm 1$	$\delta_4 = \mp 1$	$\delta_4 = -1$
$M_4'$	$\delta_2 \delta_4 = \pm 1$	$\delta_4 = \mp 1$	$\delta_3 = -1$
$M_5$	$\delta_4 = -1$	$\delta_4 = -1$	$\delta_2 = -1$
$M_5'$	$\delta_3 = -1$	$\delta_4 = -1$	$\delta_2 = -1$

By looking at Table 4 we immediately see that  $M_1$  has no pin<sup>+</sup> structures,  $M'_2$  and  $M_3$  admit no pin<sup>±</sup> structures and  $M'_3$  has no pin<sup>-</sup> structures, since the corresponding equations are not compatible. We now list all the characters  $\varepsilon_{\pm|\Lambda}$ , corresponding to the pin<sup>±</sup> and spin structures in the remaining cases.

$$\varepsilon_{-}(M_{1}) = (\delta_{1}, -\delta_{1}, -1, \delta_{4}), \qquad \varepsilon_{\pm}(M'_{1}) = (\mp 1, \pm 1, \delta_{3}, -1), 
\varepsilon_{\pm}(M_{2}) = (\delta_{1}, -1, \pm 1, \mp 1), \qquad \varepsilon_{-}(M_{3}) = (\delta_{1}, -1, \delta_{3}, -1), 
\varepsilon_{+}(M_{4}) = (\delta_{1}, \delta_{1}, \delta_{3}, -1), \qquad \varepsilon_{\pm}(M'_{4}) = (\delta_{1}, -1, -1, \mp 1), 
\varepsilon(M_{5}) = (\delta_{1}, -1, \delta_{3}, -1), \qquad \varepsilon(M'_{5}) = (-1, -1, -1, \delta_{4}).$$

Now, for each choice of  $\varepsilon_{\pm|\Lambda}$  there are  $2^2=4$  structures corresponding to the possible choices of  $\sigma_1, \sigma_2$ , hence it is easy to verify that the number of pin<sup>+</sup>, pin<sup>-</sup> or spin structures is as indicated in the theorem.

Remark 3.3. The procedure of adding appropriate characters to  $M_1, M'_1$  to obtain orientable manifolds, with  $M_1$  admitting a spin structure while  $M'_1$  does not, can be used with the remaining pairs  $M_i, M'_i, 2 \le i \le 4$ , as well. Alternately, we can also use the method described in Remark 2.4. Indeed, consider the orientable  $\mathbb{Z}_2^2$ -manifolds  $M_{d\Gamma_1}, M_{d\Gamma'_1}$  of dimension 8 obtained by doubling the Bieberbach groups  $\Gamma_1, \Gamma'_1$  (see Table 5). The resulting manifolds now carry a Kähler structure.

 $L_{b_2}$  $B_3$  $L_{b_3}$  $L_{b_2'}$ 1/21/21/21/21 1/21/21/2-1 -1 1/21 1 1 1/21/21/2-1 -1 1/2-1 1/2

Table 5.

By comparing the Sunada numbers, we see that  $M_{\mathrm{d}\Gamma_1}$  and  $M_{\mathrm{d}\Gamma_1'}$  are isospectral. Now, we look at condition  $(\varepsilon_1)$  in (2.3). For  $\mathrm{d}\Gamma_1$  we have that  $\delta_3\delta_7=-1$ ,  $\delta_1\delta_2\delta_5\delta_6=-1$  and  $\delta_1\delta_2\delta_5\delta_6=1$ . These last two equations are clearly not compatible, hence  $M_{\mathrm{d}\Gamma_1}$  admits no spin structures. On the other hand, for  $\mathrm{d}\Gamma_1'$  we get  $\delta_2\delta_6=-1$ ,  $\delta_1\delta_4\delta_5\delta_8=-1$  and  $\delta_1\delta_2\delta_5\delta_6=1$ , hence  $\varepsilon=(\delta_1,\delta_2,\delta_3,\delta_4,-\delta_1,-\delta_2,\delta_7,-\delta_4,\sigma_1e_4e_8,\sigma_2e_3e_7)$ , thus obtaining  $2^7$  spin structures in this case.

### 4. $Pin^{\pm}$ Structures on $\mathbb{Z}_2$ -manifolds.

In this last section we study in some detail the special case of  $\mathbb{Z}_2$ -manifolds, where an explicit description of the pin<sup>±</sup> structures can be given. For each  $0 \le j, h < n$ , let as in (3.1)

$$B_{j,h} := \operatorname{diag}(\underbrace{J, \dots, J}_{i}, \underbrace{-1, \dots, -1}_{h}, \underbrace{1, \dots, 1}_{l})$$

where n = 2j + h + l,  $j + h \neq 0$  and  $l \geq 1$ . Then  $B_{j,h} \in O(n)$ ,  $B_{j,h}^2 = Id$  and  $B_{j,h} \in SO(n)$  if and only if j + h is even. Let  $\Lambda = \mathbb{Z}e_1 \oplus \cdots \oplus \mathbb{Z}e_n$  be the canonical lattice of  $\mathbb{R}^n$  and for j, h as before define the groups

(4.1) 
$$\Gamma_{j,h} := \langle B_{j,h} L_{\frac{e_n}{2}}, \Lambda \rangle.$$

We have that  $\Lambda$  is stable by  $B_{j,h}$  and  $(B_{j,h} + Id)\frac{e_n}{2} = e_n \in \Lambda \setminus (B_{j,h} + Id)\Lambda$ . Hence, by Proposition 2.1 in [**DM**], the  $\Gamma_{j,h}$  are Bieberbach groups. In this way, if  $M_{j,h} = \Gamma_{j,h} \setminus \mathbb{R}^n$ , we have a family

(4.2) 
$$\mathcal{F} = \{ M_{j,h} : 0 \le j \le \left[ \frac{n-1}{2} \right], 0 \le h < n-2j, j+h \ne 0 \}$$

of compact flat manifolds with holonomy group  $F \simeq \mathbb{Z}_2$ . The next proposition summarizes some known results on  $\mathbb{Z}_2$ -manifolds. We include a proof for completeness.

**Proposition 4.1.** The family  $\mathcal{F}$  gives a system of representatives for the diffeomorphism classes of  $\mathbb{Z}_2$ -manifolds of dimension n. Furthermore we have:

(4.3) 
$$H_1(M_{j,h}, \mathbb{Z}) \simeq \mathbb{Z}^{j+l} \oplus \mathbb{Z}_2^h.$$

For  $1 \leq p \leq n$ ,

(4.4) 
$$\beta_p(M_{j,h}) = \sum_{i=0}^{\left[\frac{p}{2}\right]} {j+h \choose 2i} {j+l \choose p-2i}.$$

If  $\beta_1(M_{j,h}) = \beta_1(M_{j',h'})$ , then  $\beta_p(M_{j,h}) = \beta_p(M_{j',h'})$  for any  $p \ge 1$ .

*Proof.* We first prove that the manifolds  $M_{j,h}$  are pairwise non homeomorphic. We now compute  $H_1(M_{j,h},\mathbb{Z}) \simeq \Gamma_{j,h}/[\Gamma_{j,h},\Gamma_{j,h}]$ . For  $\gamma = B_{j,h}L_{\frac{e_n}{2}}$ , we have

$$[\Gamma_{j,h}, \Gamma_{j,h}] = \langle [\gamma, L_{e_i}] = L_{(B-Id)e_i} : 1 \le i \le n \rangle$$
  
=  $\langle L_{e_2-e_1}, \dots, L_{e_{2i}-e_{2i-1}}, L_{2e_{2i+1}}, \dots, L_{2e_{2i+h}} \rangle.$ 

Using this information and the fact that  $\gamma^2 = L_{e_n}$  it is easy to see that

$$H_1(M_{j,h},\mathbb{Z})\simeq \mathbb{Z}^{j+l}\oplus \mathbb{Z}_2^h.$$

Thus, if  $M_{j,h}$  and  $M_{j',h'}$  are homeomorphic then h = h' and j + l = j' + l', hence j = j' as asserted.

To show that the family  $\mathcal{F}$  gives a complete system of representatives for the diffeomorphism classes of  $\mathbb{Z}_2$ -manifolds, we will use results in  $[\mathbf{Ch}]$ , p.153 (it could also be proved directly by using that any integral representation of  $\mathbb{Z}_2$  decomposes uniquely as a sum of indecomposable representations of rank  $\leq 2$  given by 1, -1 or J).

The cardinality of  $\mathcal{F}$  equals  $\left(\sum_{j=0}^{\left[\frac{n-1}{2}\right]}n-2j\right)-1$ , since we must exclude the case j=h=0 corresponding to  $B_{0,0}=Id$ . Thus we have

(4.5) 
$$\#\mathcal{F} = \left(n - \left[\frac{n-1}{2}\right]\right) \left(\left[\frac{n-1}{2}\right] + 1\right) - 1 = \begin{cases} \frac{n^2 + 2n - 4}{4} & n \text{ even} \\ \frac{n^2 + 2n - 3}{4} & n \text{ odd.} \end{cases}$$

On the other hand, if p is a prime, Charlap gives a formula for the number  $N_p$  of diffeomorphism classes of  $\mathbb{Z}_p$ -manifolds of dimension n. For p=2 this number is given by:

$$N_2 = \frac{1}{2} \left[ \frac{n-1}{2} \right] \left( \left[ \frac{n-1}{2} \right] + 3 \right) + \frac{1}{2} \left( (n-1) - \left[ \frac{n-1}{2} \right] \right) \left( n - \left[ \frac{n-1}{2} \right] \right).$$

In this way we obtain that  $N_2 = \frac{1}{8}(n-2)(n+4) + \frac{1}{8}n(n+2) = \frac{n^2+2n-4}{4}$  for n even, and  $N_2 = \frac{1}{4}(n-1)(n+3) = \frac{n^2+2n-3}{4}$ , for n odd. This shows that  $\#\mathcal{F} = N_2$ , as claimed.

To determine the p-Betti number of  $M_{j,h}$  for  $1 \leq p \leq n$ , we note that  $B_{j,h}$  acts diagonally on the basis  $e_1 \pm e_2, \ldots, e_{2j-1} \pm e_{2j}, e_{2j+1}, \ldots, e_n$ , with j+l (resp. j+h) eigenvectors with eigenvalue 1 (resp. -1). Thus, an exterior product of p elements of this basis will be invariant by  $B_{j,h}$ , if and only if an even number of them have eigenvalue -1. Hence we have

$$\beta_p(M_{j,h}) = \sum_{i=0}^{\left[\frac{p}{2}\right]} \binom{j+h}{2i} \binom{j+l}{p-2i}$$

as asserted. Now, if  $\beta_1(M_{j,h}) = \beta_1(M_{j',h'})$  then j+l=j'+l' and hence j+h=j'+h'. Thus,  $\beta_p(M_{j,h}) = \beta_p(M_{j',h'})$ , for any  $1 \leq p \leq n$ .

The next result gives a description of  $pin^{\pm}$  structures on  $\mathbb{Z}_2$ -manifolds.

**Proposition 4.2.** Every  $\mathbb{Z}_2$ -manifold  $M_{\Gamma}$  has  $pin^{\pm}$  structures (and spin structures, if  $M_{\Gamma}$  is orientable). If  $\Gamma = \Gamma_{j,h}$  then  $M_{\Gamma}$  has  $2^{n-j}$   $pin^{\pm}$  structures parametrized by the tuples  $(\delta_1, \ldots, \delta_n, \sigma) \in \{\pm 1\}^{n+1}$  satisfying:

$$\delta_1 = \delta_2, \cdots, \delta_{2j-1} = \delta_{2j}$$

and

(4.7) 
$$\delta_n = \begin{cases} (-1)^{jh} (-1)^{\left[\frac{j}{2}\right]} (-1)^{\left[\frac{h}{2}\right]} & \text{for } pin^+ \text{ structures} \\ (-1)^{jh} (-1)^{\left[\frac{j+1}{2}\right]} (-1)^{\left[\frac{h+1}{2}\right]} & \text{for } pin^- \text{ structures.} \end{cases}$$

In particular, in the case of spin structures we have  $\delta_n = (-1)^{\frac{j+h}{2}}$ .

*Proof.* In light of Proposition 4.1, we have that  $\Gamma \simeq \Gamma_{j,h}$  for some j,h, hence  $M_{\Gamma}$  is diffeomorphic to  $M_{j,h}$ . Therefore, since pin<sup>±</sup> structures on diffeomorphic manifolds are in a bijective correspondence, we may assume that  $\Gamma = \Gamma_{j,h}$ .

We have observed in Remark 2.2 that equation  $(\varepsilon_2)$  always holds for  $\mathbb{Z}_2$ -manifolds of diagonal type. However, in the non diagonal case,  $(\varepsilon_2)$  gives a restriction. Namely, let  $\lambda = \sum_{i=1}^{n} m_i e_i$ ,  $m_i \in \mathbb{Z}$ . Then

$$(B_{j,h} - Id)\lambda = \sum_{i=1}^{j} (m_{2i} - m_{2i-1})e_{2i-1} + (m_{2i-1} - m_{2i})e_{2i} - 2\sum_{i=1}^{h} m_{2j+i}e_{2j+i}.$$

Thus,  $(\varepsilon_2)$  holds if and only if

$$\delta_1^{(m_2-m_1)}\delta_2^{(m_1-m_2)}\cdots\delta_{2j-1}^{(m_{2j}-m_{2j-1})}\delta_{2j}^{(m_{2j-1}-m_{2j})}=1$$

for every  $m_1, \ldots, m_{2j} \in \mathbb{Z}$ , or equivalently,

$$\delta_1 = \delta_2, \ldots, \ \delta_{2j-1} = \delta_{2j}.$$

Each of these relations divides by 2 the number of structures. Hence we obtain a maximum of  $2^{n-j+1}$  pin<sup>±</sup> structures for  $M_{j,h}$ . Furthermore, equation  $(\varepsilon_1)$  gives another restriction since  $\varepsilon_{\pm}(\gamma^2) = \varepsilon_{\pm}(L_{(B+Id)b}) = \varepsilon_{\pm}(\gamma)^2$ . Now  $(B+Id)b = e_n$ , hence, by (3.3), equation  $(\varepsilon_1)$  reads:

(4.8) 
$$\delta_n = \begin{cases} (-1)^{jh} (-1)^{\left[\frac{j}{2}\right]} (-1)^{\left[\frac{h}{2}\right]} & \text{in } Cl^+(n) \\ (-1)^{jh} (-1)^{\left[\frac{j+1}{2}\right]} (-1)^{\left[\frac{h+1}{2}\right]} & \text{in } Cl^-(n). \end{cases}$$

Thus, the restriction imposed by (4.8) divides by 2 the number of structures and we get a total of  $2^{n-j}$  pin<sup>±</sup> structures on  $M_{\Gamma}$  for  $\Gamma = \Gamma_{j,h}$ .

*Note.* Proposition 4.1 together with Lemma 3.1, give an explicit description of all pin<sup> $\pm$ </sup> structures on  $\mathbb{Z}_2$ -manifolds.

**Example 4.3.** As a final task, to illustrate Proposition 4.2, we list explicitly the 28 pin<sup> $\pm$ </sup> Riemannian  $\mathbb{Z}_2$ -manifolds  $(M, \varepsilon)$  of dimension 3 having canonical lattice of translations  $\Lambda$ .

There are 3 diffeomorphism classes, one of which splits into 2 isometry classes, hence we have 4 isometry classes, corresponding to the groups  $\Gamma_{1,0} = \left\langle \begin{bmatrix} J_1 \end{bmatrix} L_{\frac{e_3}{2}}, \Lambda \right\rangle$ ,  $\Gamma_{0,1} = \left\langle \begin{bmatrix} -1 \\ 1 \end{bmatrix} L_{\frac{e_3}{2}}, \Lambda \right\rangle$ ,  $\Gamma_{0,1} = \left\langle \begin{bmatrix} -1 \\ 1 \end{bmatrix} L_{\frac{e_3}{2}}, \Lambda \right\rangle$  and  $\Gamma_{0,2} = \left\langle \begin{bmatrix} -1 \\ -1 \end{bmatrix} L_{\frac{e_3}{2}}, \Lambda \right\rangle$ .

We note that  $M_{0,1}$  and  $M'_{0,1}$  are not isometric, as can be seen by computing the injectivity radius, that is the length of the shortest closed geodesic. Indeed, using the results in [MR3] one easily sees that these equal  $\frac{1}{2}$  and  $\frac{\sqrt{2}}{2}$ , respectively.

The pairs  $(M, \varepsilon)$  are listed in the following table, obtained by using Lemma 3.1 and Theorem 2.1.

Table 6.  $Pin^{\pm}$  structures on  $\mathbb{Z}_2$ -manifolds of dimension 3.

$M_{\Gamma}$	cond. $(\varepsilon_1)$	cond. $(\varepsilon_2)$	pin <sup>±</sup> structures	#
$M_{1,0}$	$\delta_3 = \pm 1$	$\delta_1 = \delta_2$	$\varepsilon_{\pm} = (\delta_1, \delta_1, \pm 1; \sigma \frac{\sqrt{2}}{2} (e_1 - e_2))$	$2^2$
$M_{0,1}$	$\delta_3 = \pm 1$	_	$\varepsilon_{\pm} = (\delta_1, \delta_2, \pm 1; \sigma e_1)$	$2^3$
$M'_{0,1}$	$\delta_2 \delta_3 = \pm 1$		$\varepsilon_{\pm} = (\delta_1, \delta_2, \pm \delta_2; \sigma e_1)$	$2^3$
$M_{0,2}$	$\delta_3 = -1$	_	$\varepsilon_{\pm} = (\delta_1, \delta_2, -1; \sigma e_1 e_2)$	$2^3$

We note that the spin structures for  $M_{0,2}$  are already contained in [Pf].

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